

Identification of Vacancy Defects using Positron Annihilation

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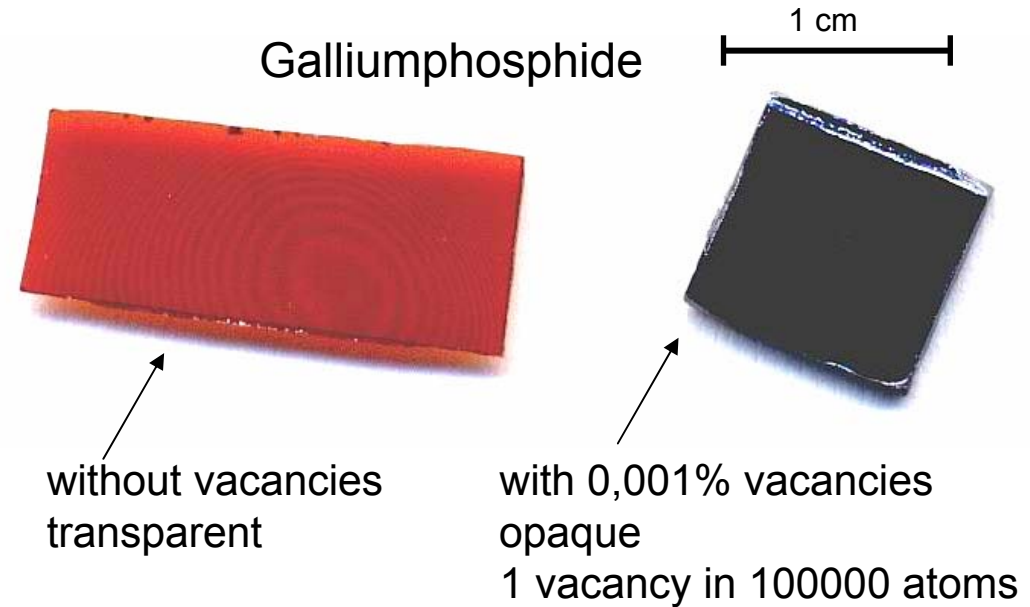
Martin-Luther-Universität
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- Why are so tiny defects important at all?
- Positron trapping by vacancies
- Other techniques of positron annihilation
- Vacancy clusters



Point defects determine optical and electronic properties of semiconductors

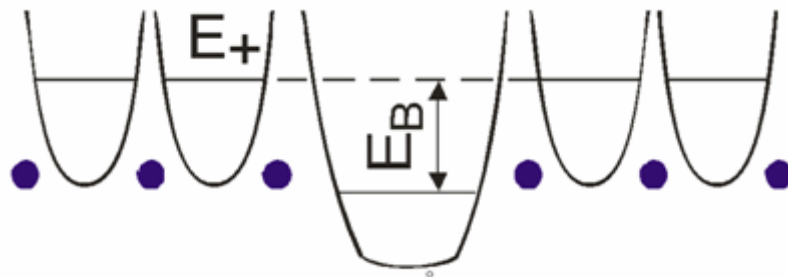
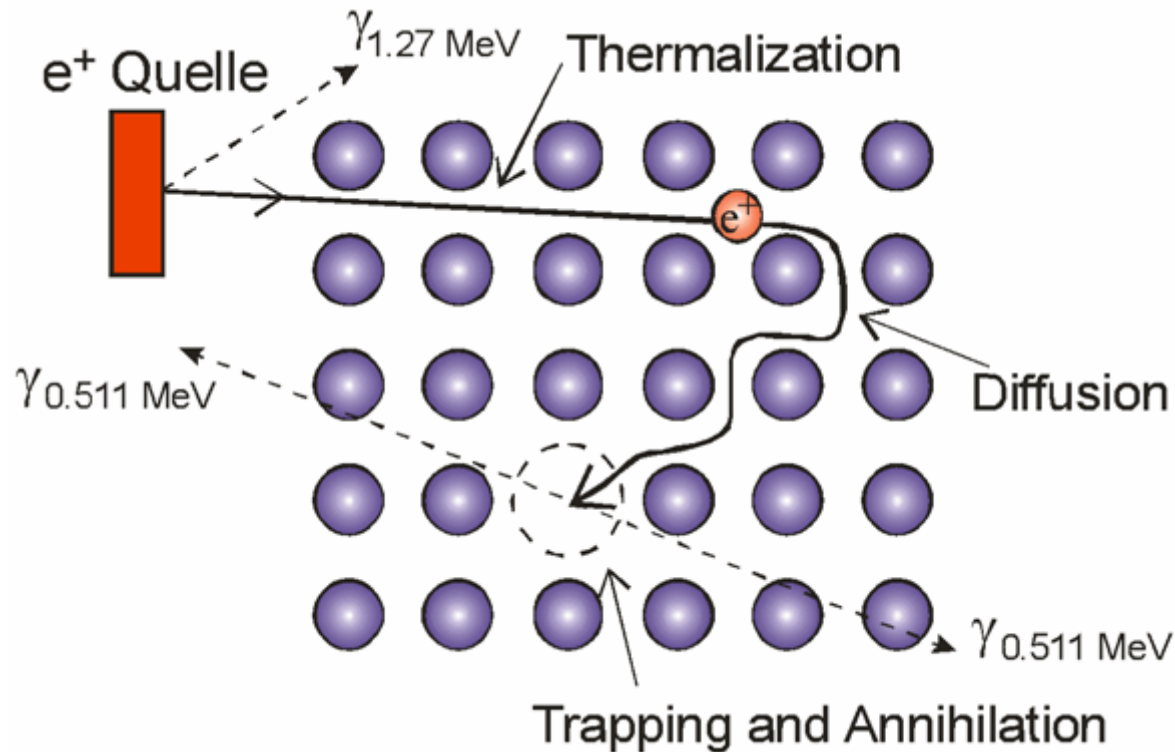
- Point defects determine electronic and optical properties
- electric conductivity strongly influenced
- Doping of semiconductors (n-, p-Si)



- Point defects are generated by irradiation (e.g. cosmic rays), by plastic deformation or by diffusion, ...
- Metals in high radiation environment -> formation of voids -> embrittlement
- -> Properties of vacancies and other point defects must be known
- Analytical tools are needed to characterize point defects

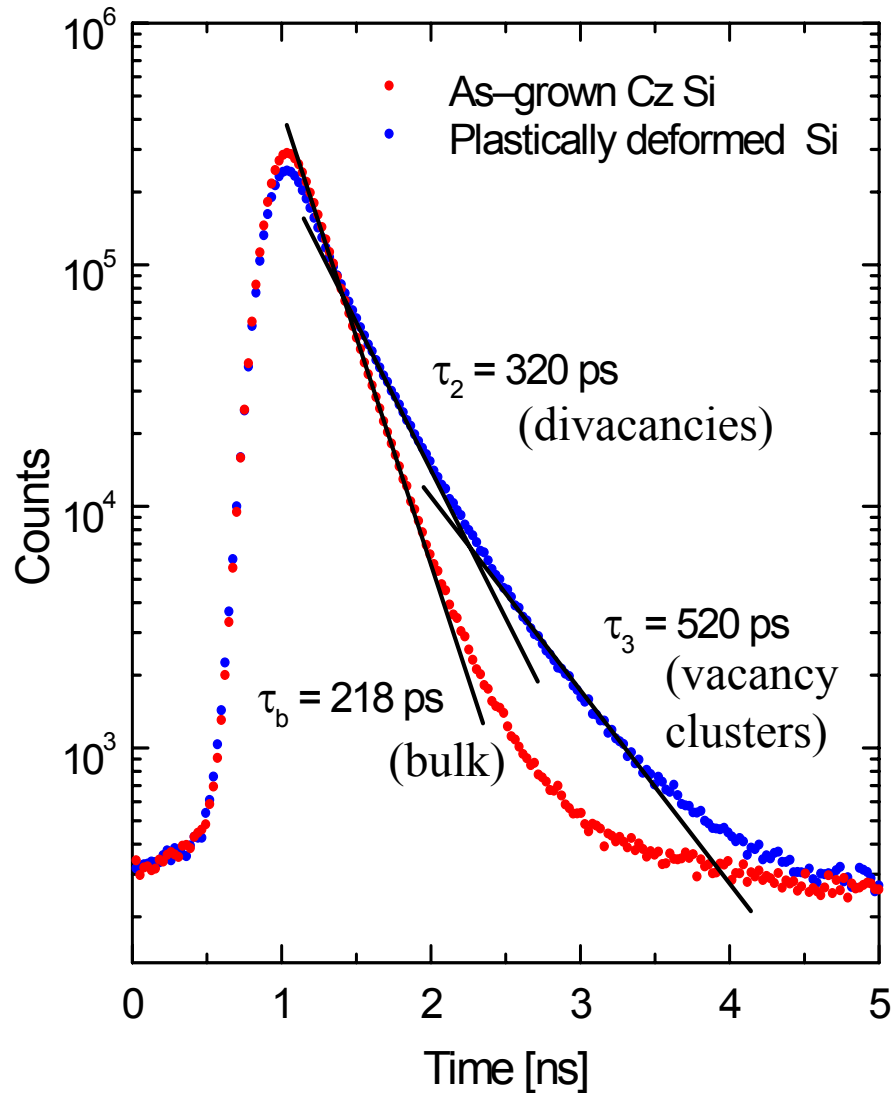
The positron lifetime spectroscopy

^{22}Na



- positron wave-function can be localized in the attractive potential of a defect
- annihilation parameters change in the localized state
- e.g. positron lifetime increases in a vacancy
- lifetime is measured as time difference between appearance of 1.27 (start) and 0.51 MeV (stop) quanta
- defect identification and quantification possible

Positron lifetime spectroscopy



- positron lifetime spectra consist of exponential decay components
- positron trapping in open-volume defects leads to long-lived components
- longer lifetime due to lower electron density
- analysis by non-linear fitting: lifetimes τ_i and intensities I_i

positron lifetime spectrum:

$$N(t) = \sum_{i=1}^{k+1} \frac{I_i}{\tau_i} \exp\left(-\frac{t}{\tau_i}\right)$$

trapping coefficient

$$K_d = \mu C_d = \frac{I_2}{I_1} \left(\frac{1}{\tau_b} - \frac{1}{\tau_d} \right)$$

trapping rate

defect concentration



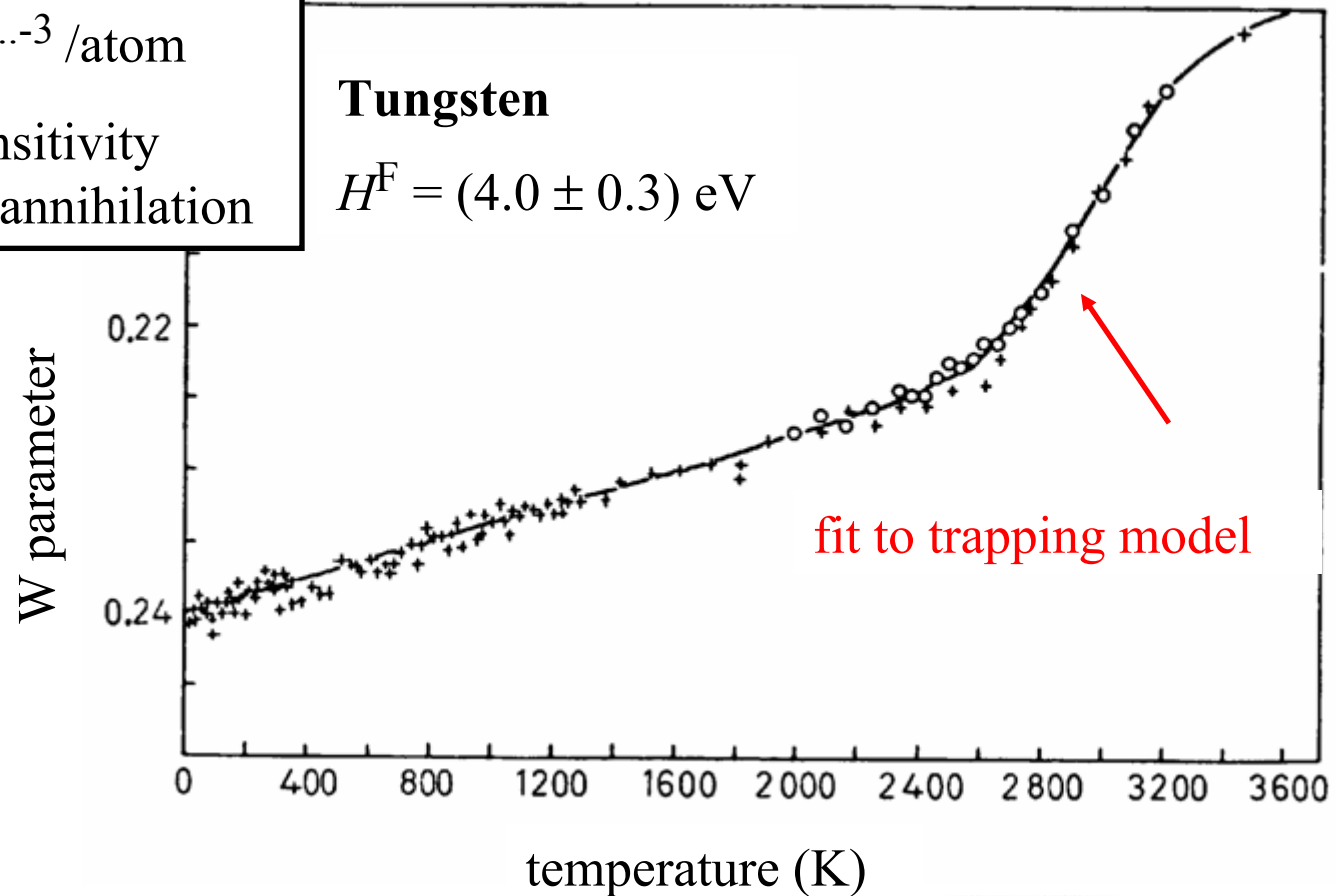
Vacancies in thermal Equilibrium

- Vacancy concentration in thermal equilibrium:
- in metals $H^F \approx 1...4 \text{ eV} \Rightarrow$ at $T_m [1v] \approx 10^{-4...-3} / \text{atom}$
- fits well to the sensitivity range of positron annihilation

$$C_{1v}(T) = \exp\left(\frac{S_{1v}^F}{k}\right) \exp\left(\frac{H_{1v}^F}{kT}\right)$$

Tungsten

$$H^F = (4.0 \pm 0.3) \text{ eV}$$



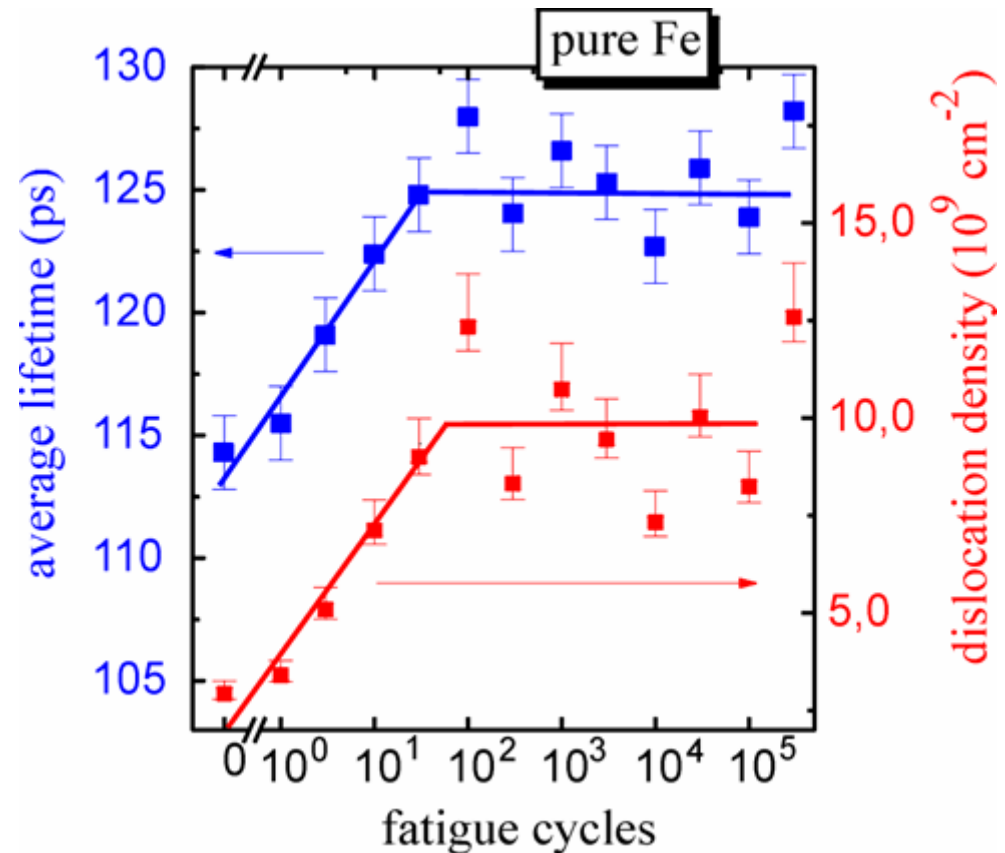
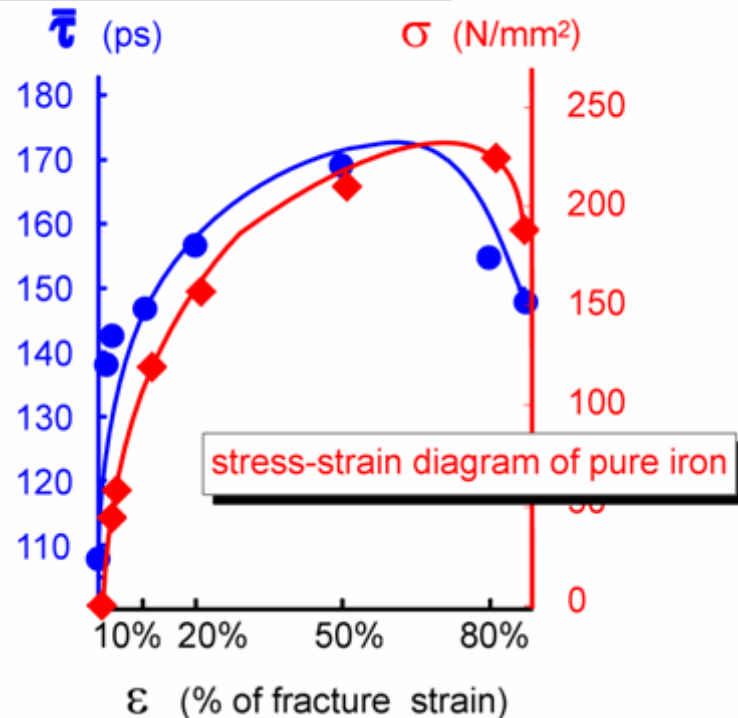
(Ziegler, 1979)



Defects in Iron after tensile strength and fatigue treatment

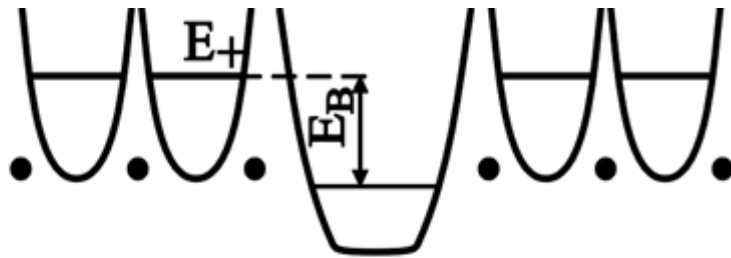
- We performed an extensive study of defects in mechanically damaged iron and steel
- Positrons are very sensitive: detection of defects already in the Hooks range of the stress-strain experiment
- Vacancy cluster and dislocations are detectable in both cases

average positron lifetime in pure iron after tensile strain

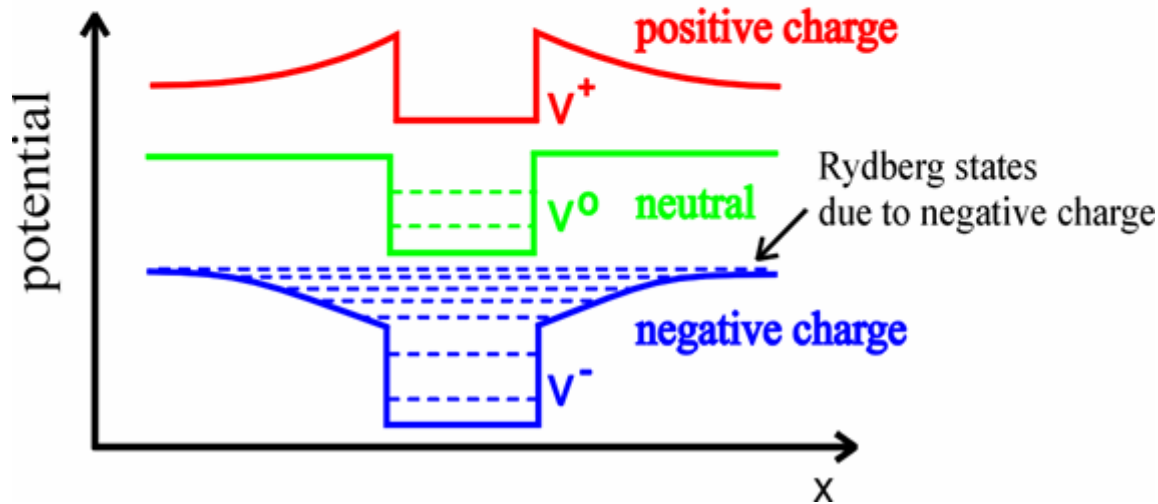


Somieski et al., J. Physique IV 5, C1/127-134 (1995)



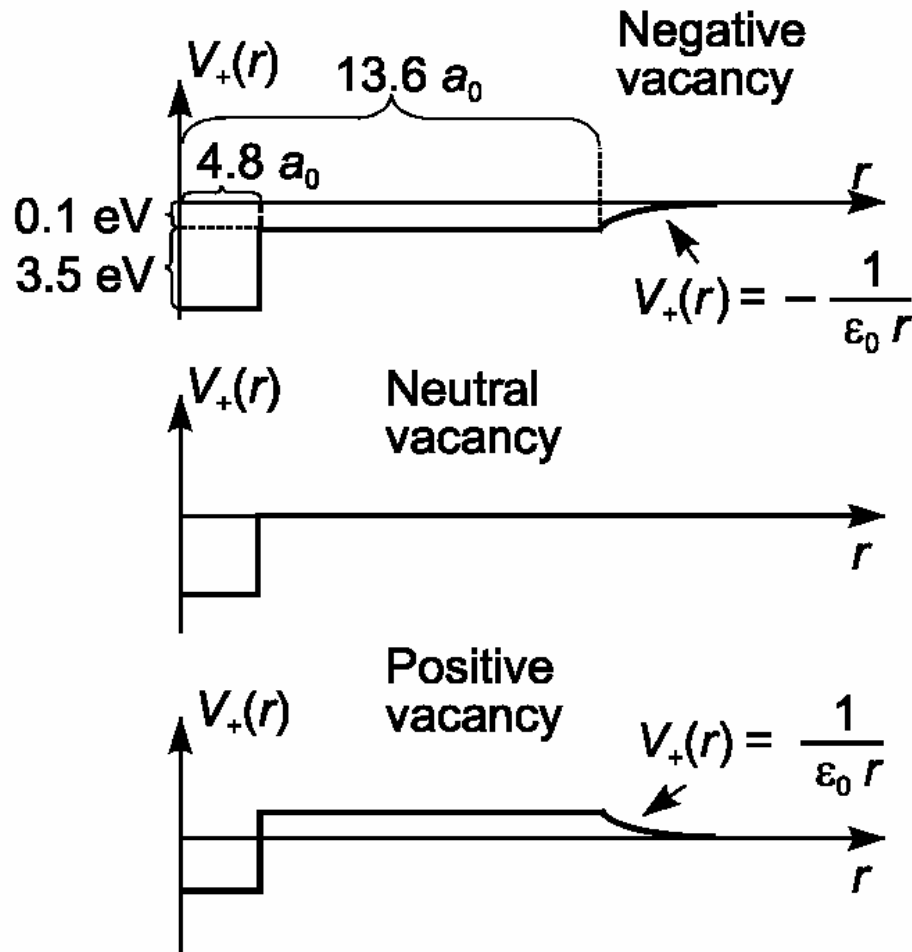


Vacancies in a semiconductor may be charged



- in a metal: charge of a vacancy is effectively screened by free electrons
- they are not available in semiconductors
- thus, long-range Coulomb potential added
- positrons may be attracted or repelled
- trapping coefficient μ is function of charge state

Vacancies may be charged



Puska et al. 1990

For a negative vacancy:

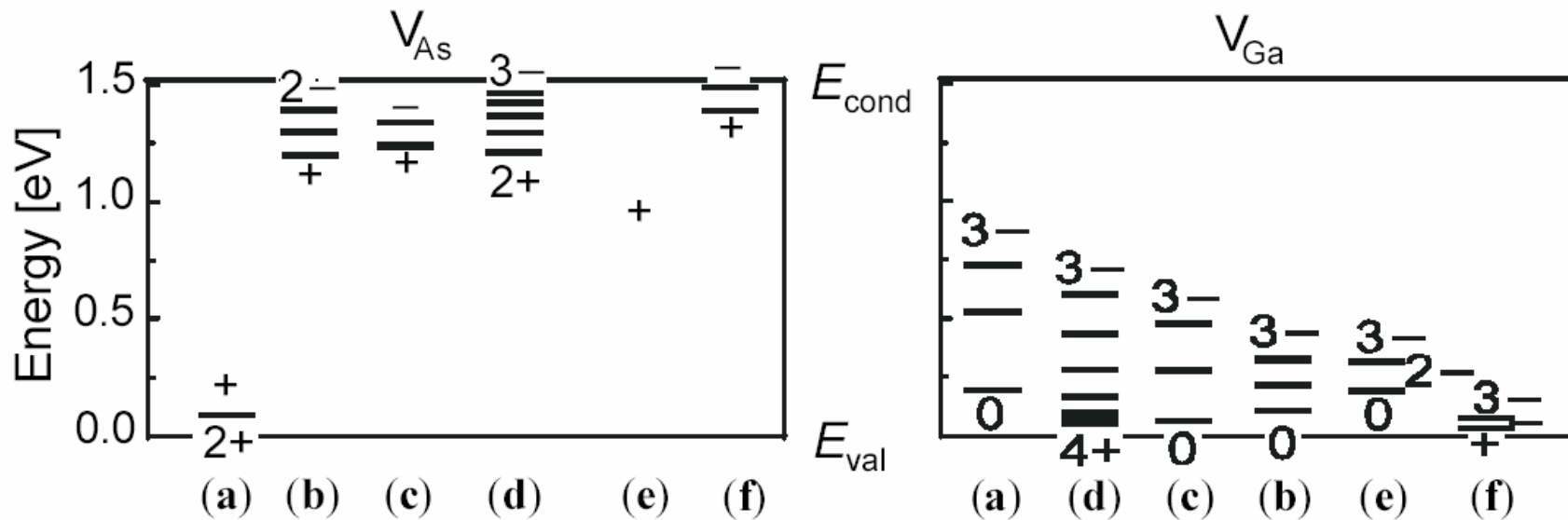
- Coulomb potential is rather extended but weak
- it supports trapping only at low temperatures
- at higher temperatures: detrapping dominates and vacancy behaves like a vacancy in a metal or a neutral vacancy

Positive vacancies repel positrons



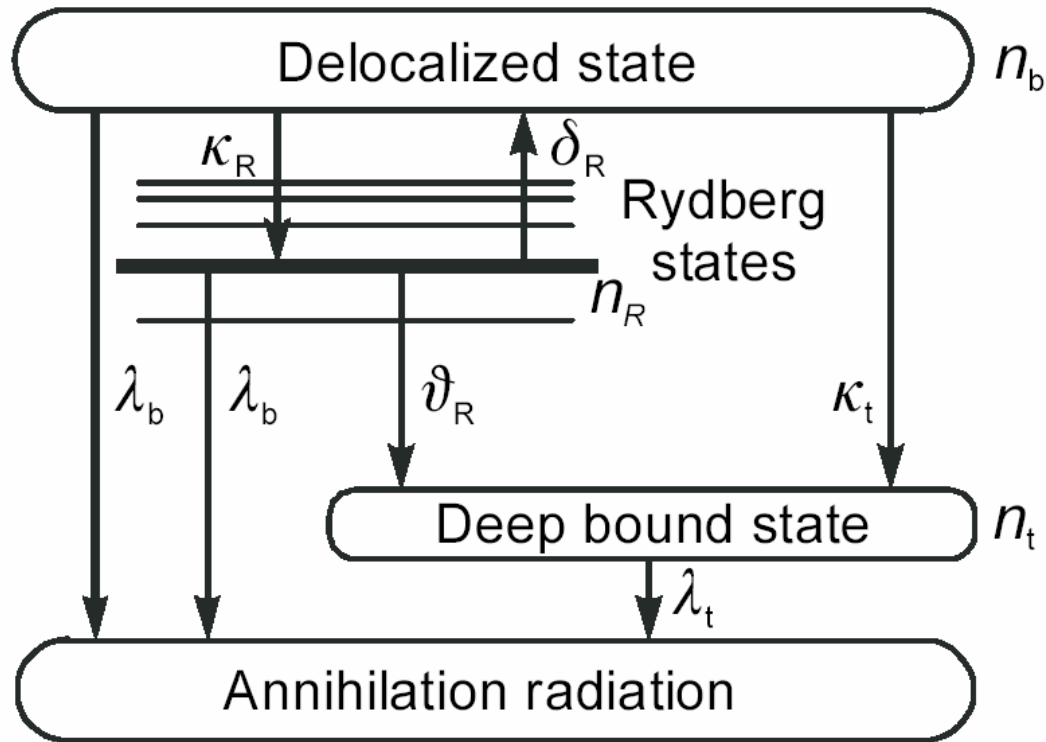
Theoretical calculation of vacancy levels in GaAs

- Theoretical description not simple
- relaxation of vacancy possible → Jahn-Teller distortion / negative-U effect



Ionization levels of arsenic vacancies, gallium vacancies, and antisites according to theoretical calculations of (a) Baraff and Schlüter (1985a), (b) Puska (1989a), (c) Jansen and Sankey (1989), (d) Xu and Lindefelt (1990), (e) Zhang and Northrup (1991), and (f) Seong and Lewis (1995), (g) Zhang and Chadi (1990), (h) Pöykkö et al. (1997). E_{val} and E_{cond} are the edges of the valence and the conduction band, respectively.

Positron trapping by negative vacancies



- trapping process can be described quantitatively by trapping model
- Coulomb potential leads to Rydberg states
- from there: positrons may re-escape by thermal stimulation
- once in the deep state: positron is captured until annihilation
- detrapping is strongly temperature dependent

$$\delta_R = \frac{\kappa_R}{\rho_v} \left(\frac{m^* k_B T}{2\pi \hbar^2} \right)^{3/2} \exp\left(-\frac{E_R}{k_B T} \right)$$

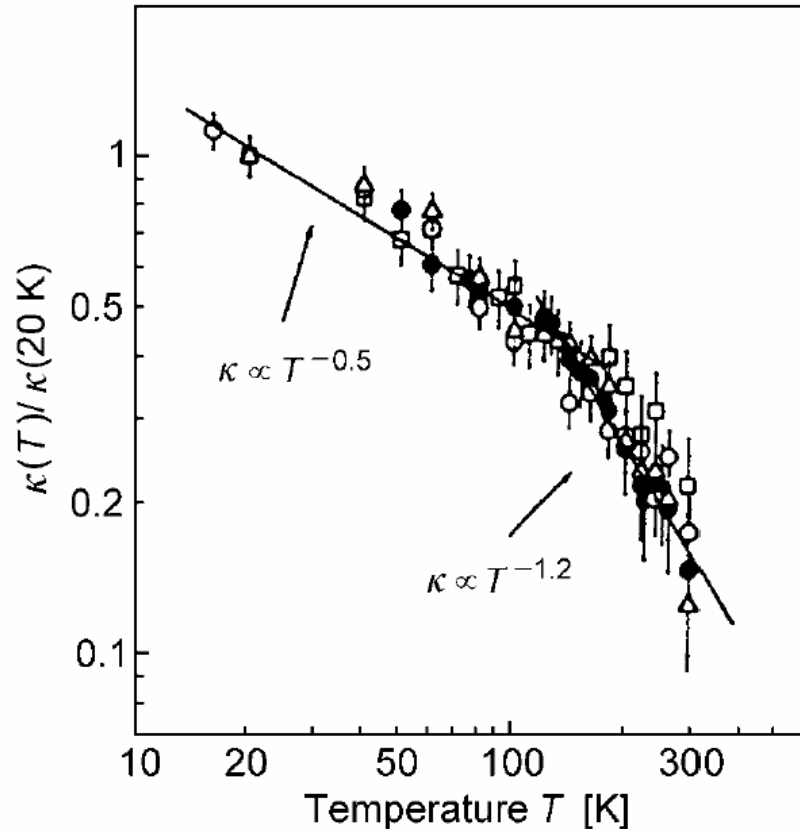
E_R binding energy of positron in Rydberg state

ρ_v vacancy density

Manninen, Nieminen, 1981



Negative vacancies show temperature-dependent positron trapping



positron trapping in negatively charged Ga vacancies in Si-GaAs

Le Berre et al., 1995

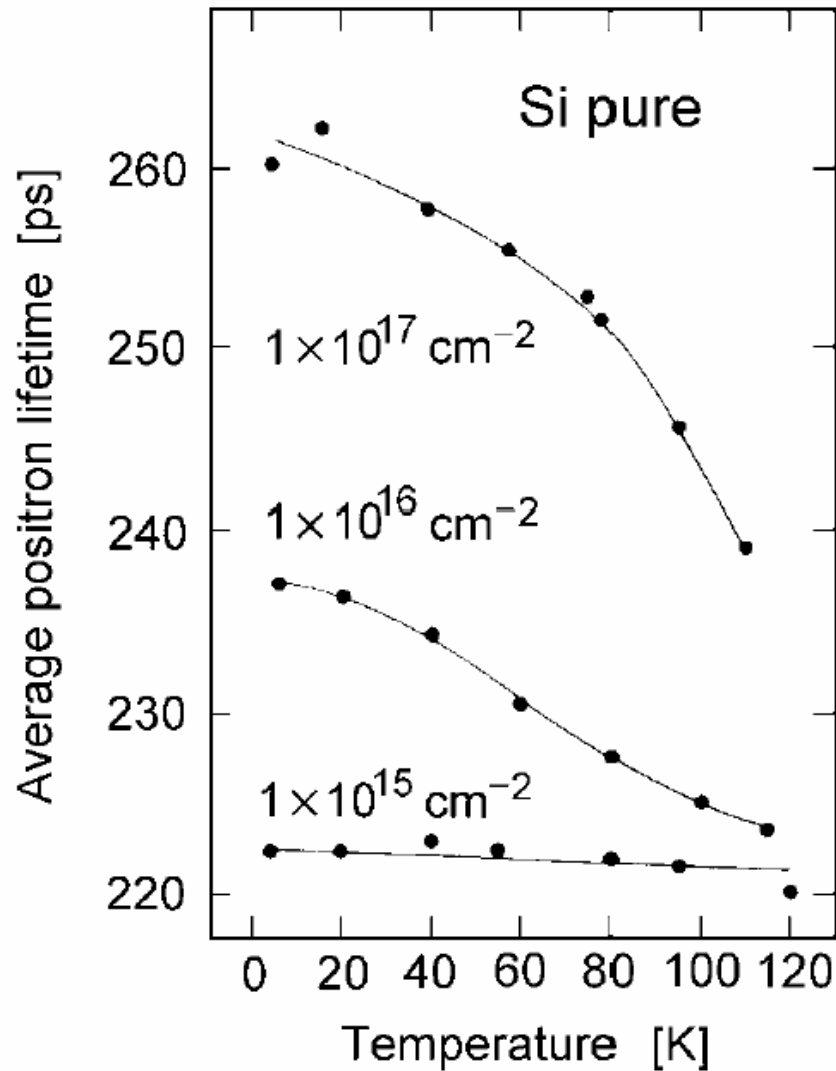
- temperature dependence of positron trapping is rather complex

$$\kappa = \frac{\vartheta_R \rho_v \kappa_{R0} T^{-1/2}}{\vartheta_R \rho_v + \kappa_{R0} \left(\frac{m^* k_B}{2\pi \hbar^2} \right)^{3/2} T \exp\left(-\frac{E_R}{k_B T}\right)}$$

- low temperature: $\sim T^{-0.5}$ due to diffusion limitation in Rydberg states
- higher T: stronger temperature dependence due to thermal detrapping from Rydberg state



Temperature-dependent positron trapping

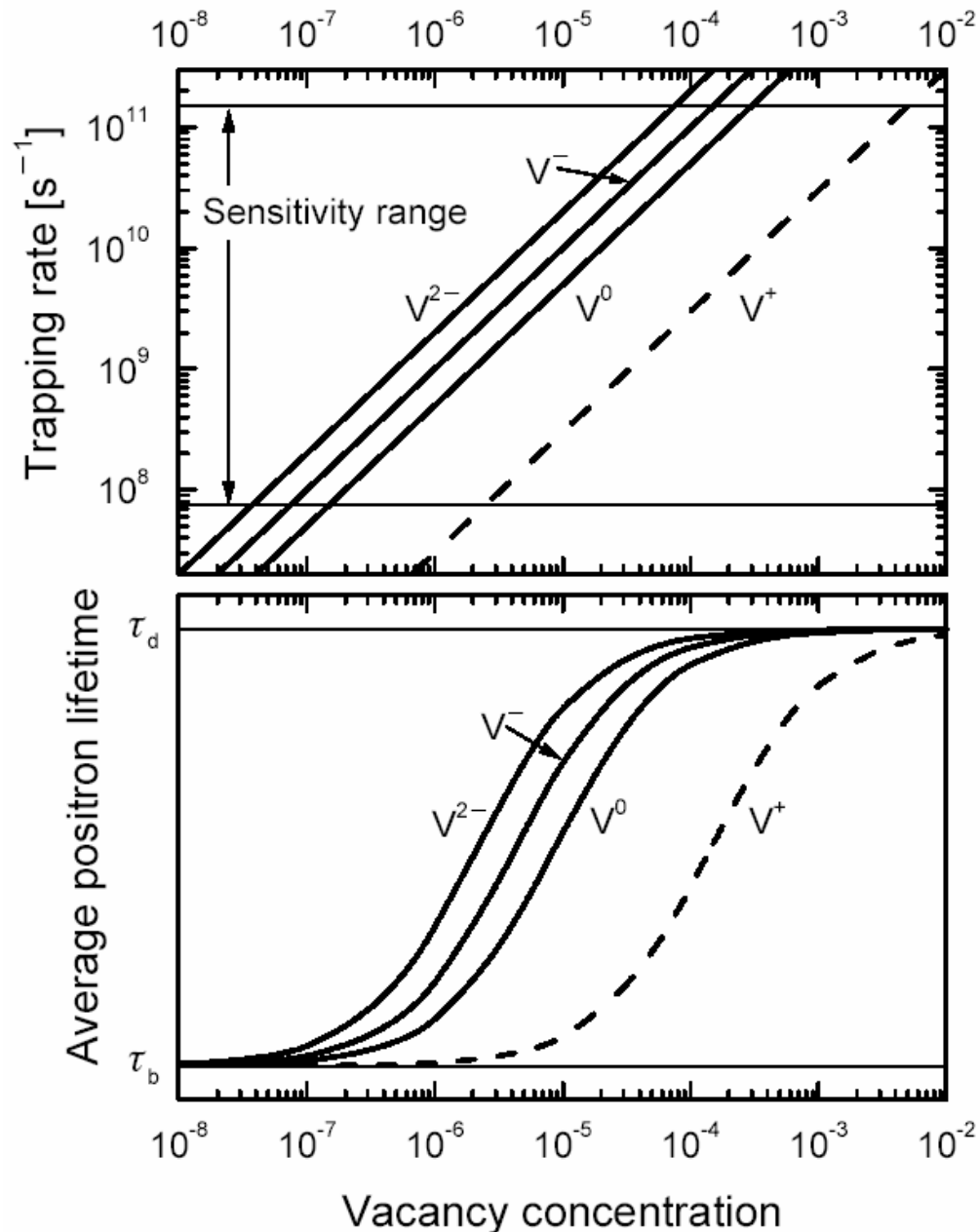


- temperature dependence of positron trapping can be used to determine the charge state of vacancies
- trapping to positive vacancies possible at elevated T
- however: has never been observed
- example: Positron trapping in e-irradiated Si
- trapping by negatively charged divacancies

(Mäkinen et al. 1989)



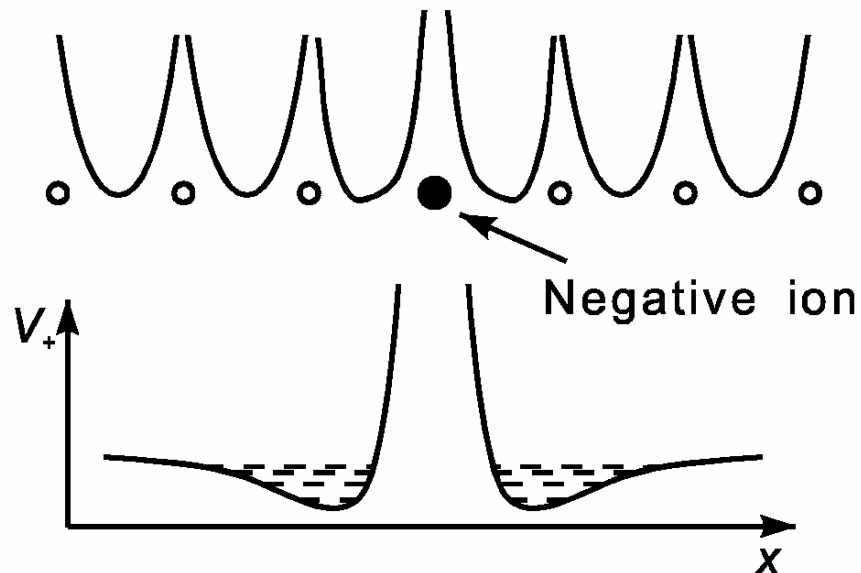
Sensitivity limits of PAS for vacancy detection



- **lower sensitivity limit** e.g. for negatively charged divacancies in Si starts at about 10^{15} cm^{-3}
- **upper limit**: saturated positron trapping
- defect identification still possible
- Then: only lower limit for defect density can be given



Negative ions act as shallow positron traps



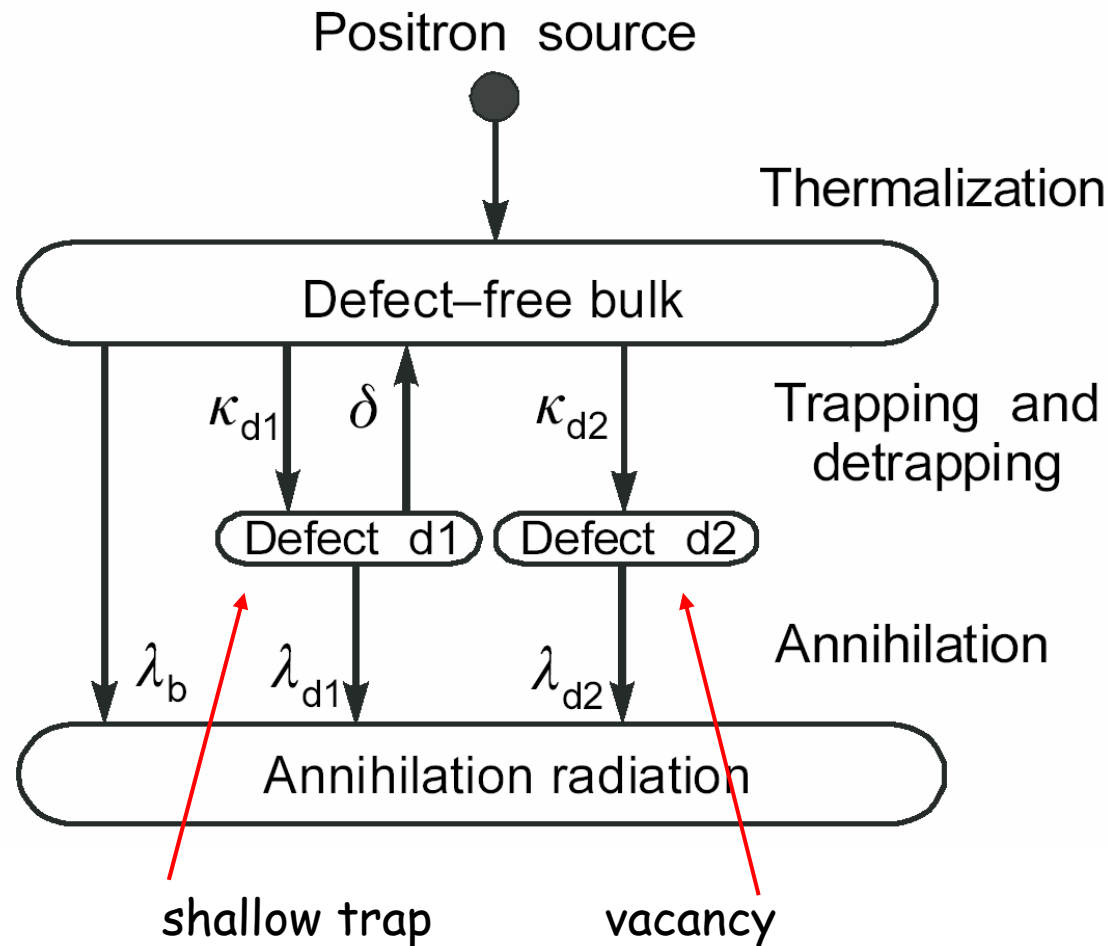
- at low T: negatively charged defects without open volume may trap positrons
- “shallow” due to small positron binding energy
- annihilation parameters close to bulk parameters
- acceptor-type impurities, dopants, negative antisite defects
- thermally stimulated detrapping can be described by:

$$\delta = \frac{\kappa}{\rho_{\text{st}}} \left(\frac{m^* k_{\text{B}} T}{2\pi\hbar^2} \right)^{3/2} \exp\left(-\frac{E_{\text{st}}}{k_{\text{B}} T} \right)$$

Saarinen et al., 1989



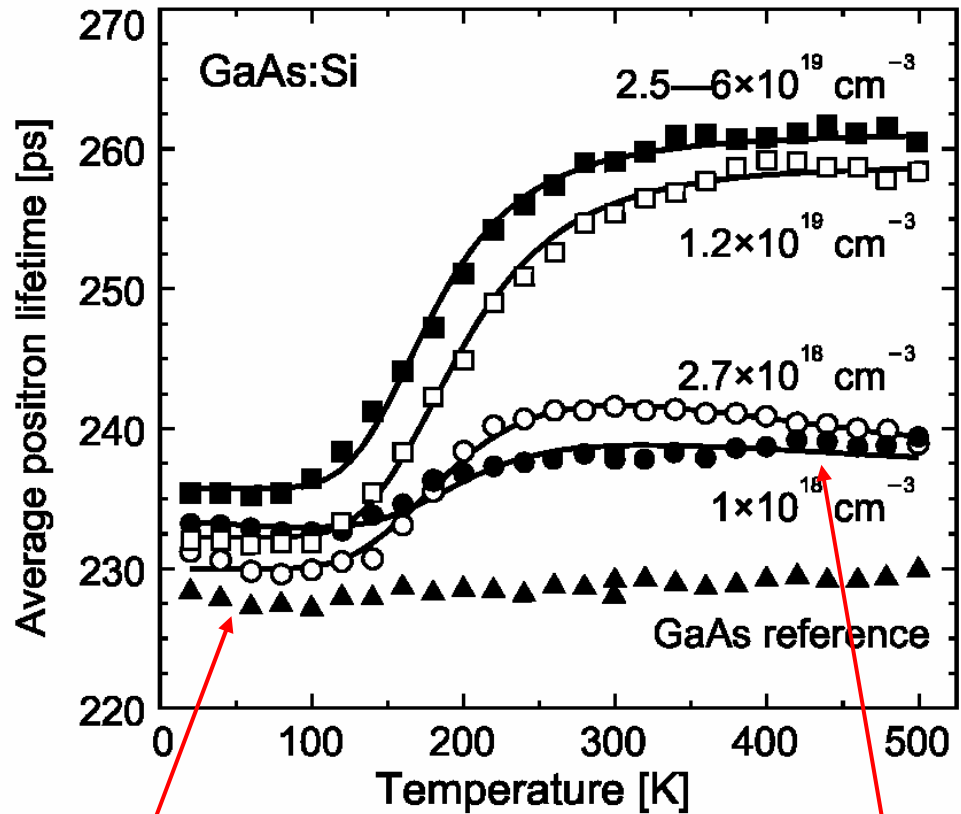
Shallow positron traps



- positron trapping model gets more complex
- however: trapping at shallow traps can be avoided at high temperatures



Effect of shallow positron traps



- temperature dependence is characterized by competing trapping by vacancies and shallow traps
- in GaAs:Si we observe $V_{Ga}-Si_{Ga}$ complexes at high temperatures
- and Si_{Ga}^- donors at low T in addition (shallow traps)

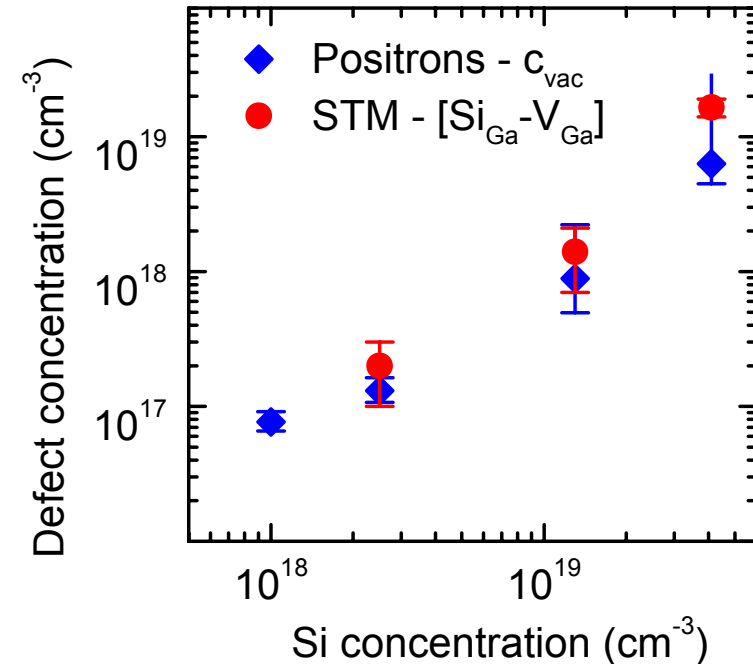
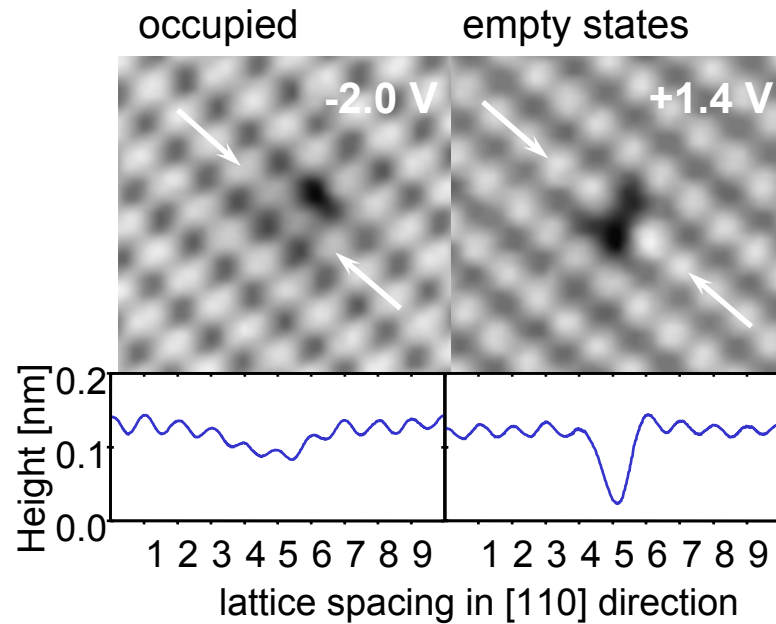
competing trapping centres at low T
shallow positron traps (Si_{Ga}^-)

trapping by vacancies
at elevated temperatures ($V_{Ga}-Si_{Ga}$)

J. Gebauer et al. 1997



Identification of $V_{Ga}-Si_{Ga}$ -Complexes in GaAs:Si



- Scanning tunneling microscopy at GaAs (110)-cleavages planes (by Ph. Ebert, Jülich)
- Defect complex identified as $V_{Ga}-Si_{Ga}$

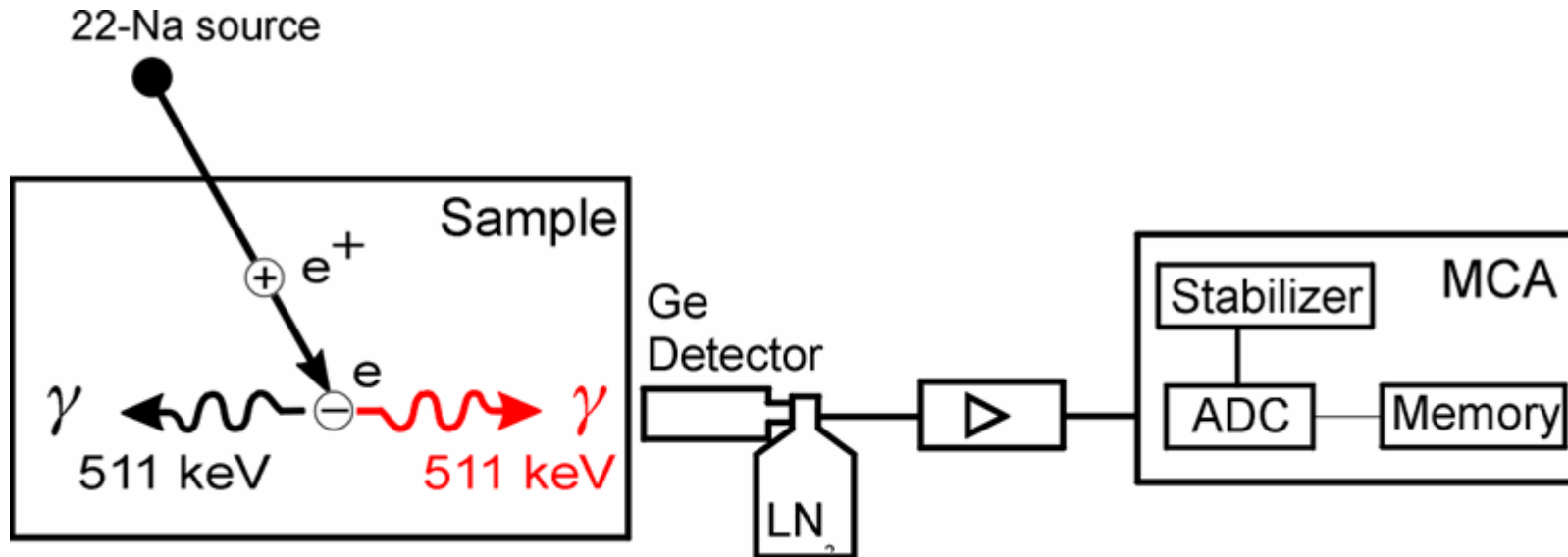
- Quantification → Agreement

Mono-vacancies in GaAs:Si are $V_{Ga}-Si_{Ga}$ -complexes

Gebauer et al., Phys. Rev. Lett. **78** (1997) 3334

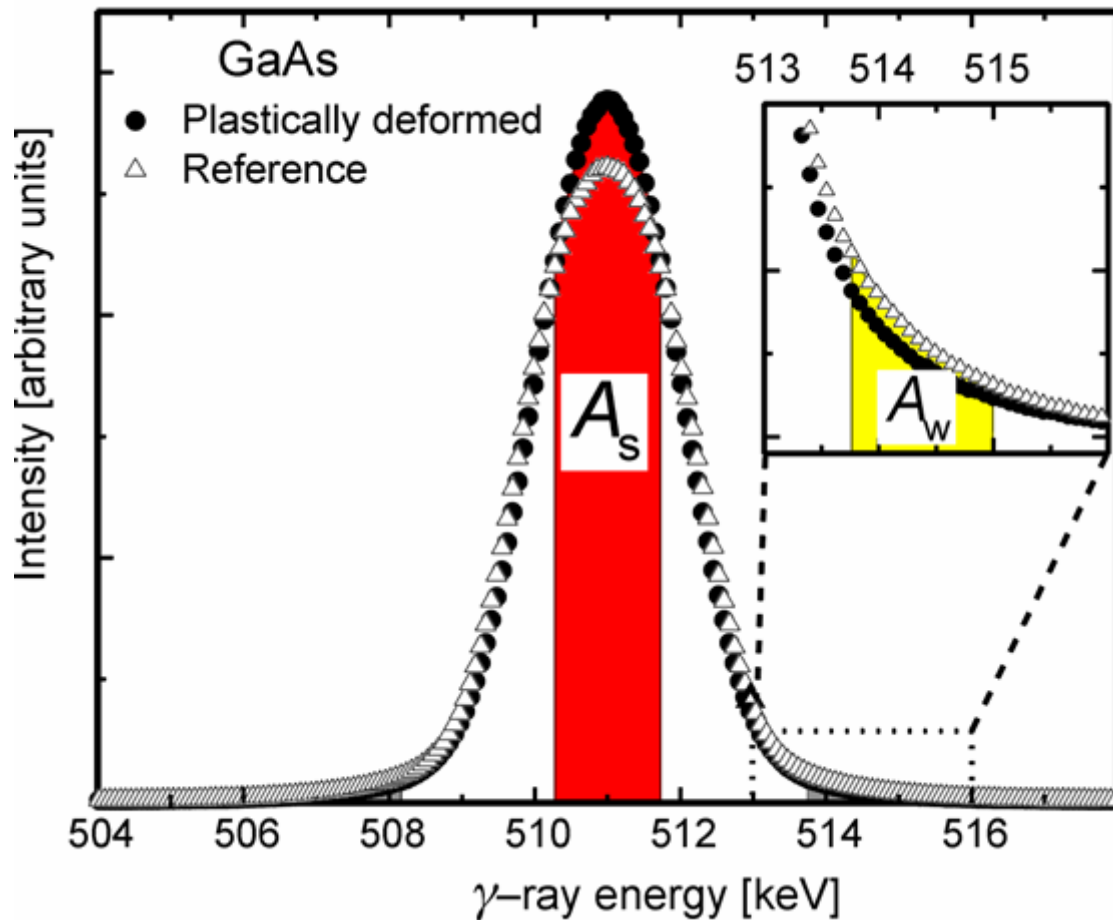


Measurement of Doppler Broadening



- electron momentum in propagation direction of 511 keV γ -ray leads to Doppler broadening of annihilation line
- can be detected by conventional energy-dispersive Ge detectors and standard electronics

Line Shape Parameters



S parameter:

$$S = A_S/A_0$$

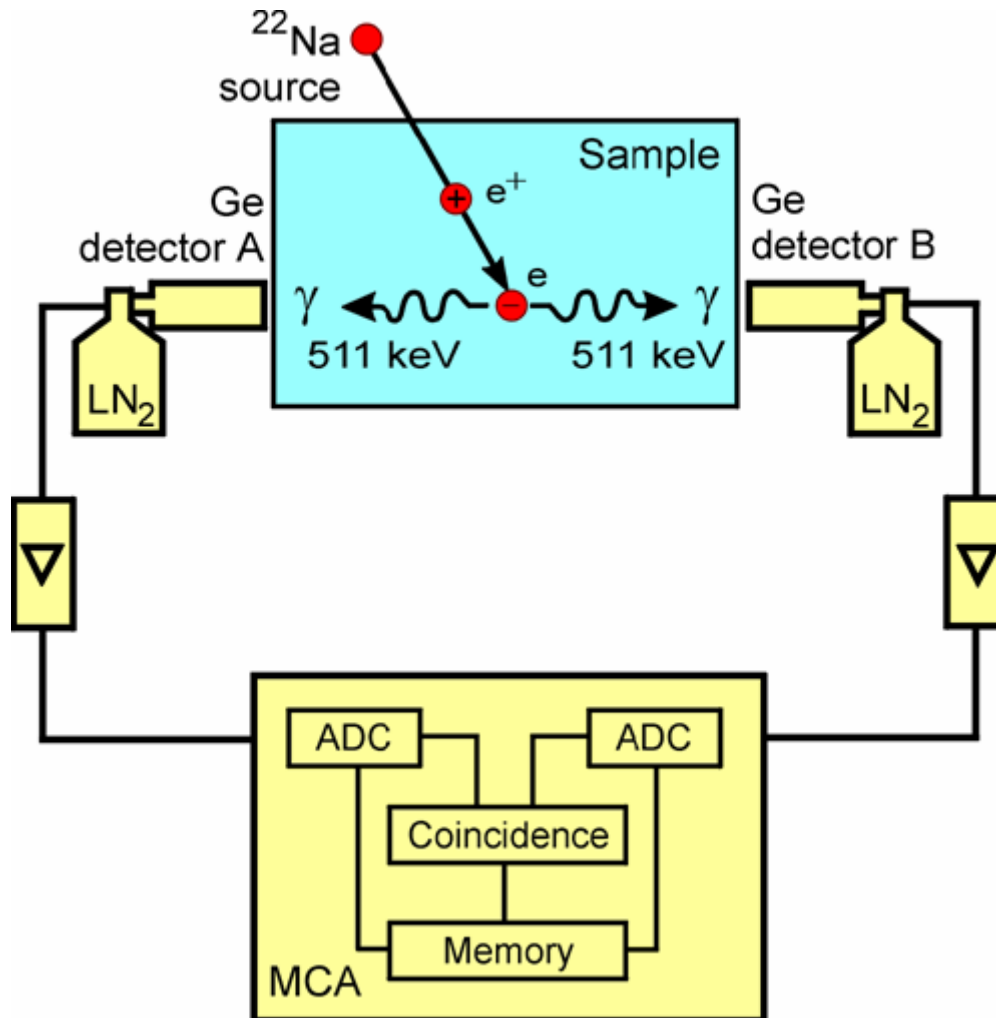
W parameter:

$$W = A_W/A_0$$

W parameter mainly determined by annihilations of core electrons (chemical information)

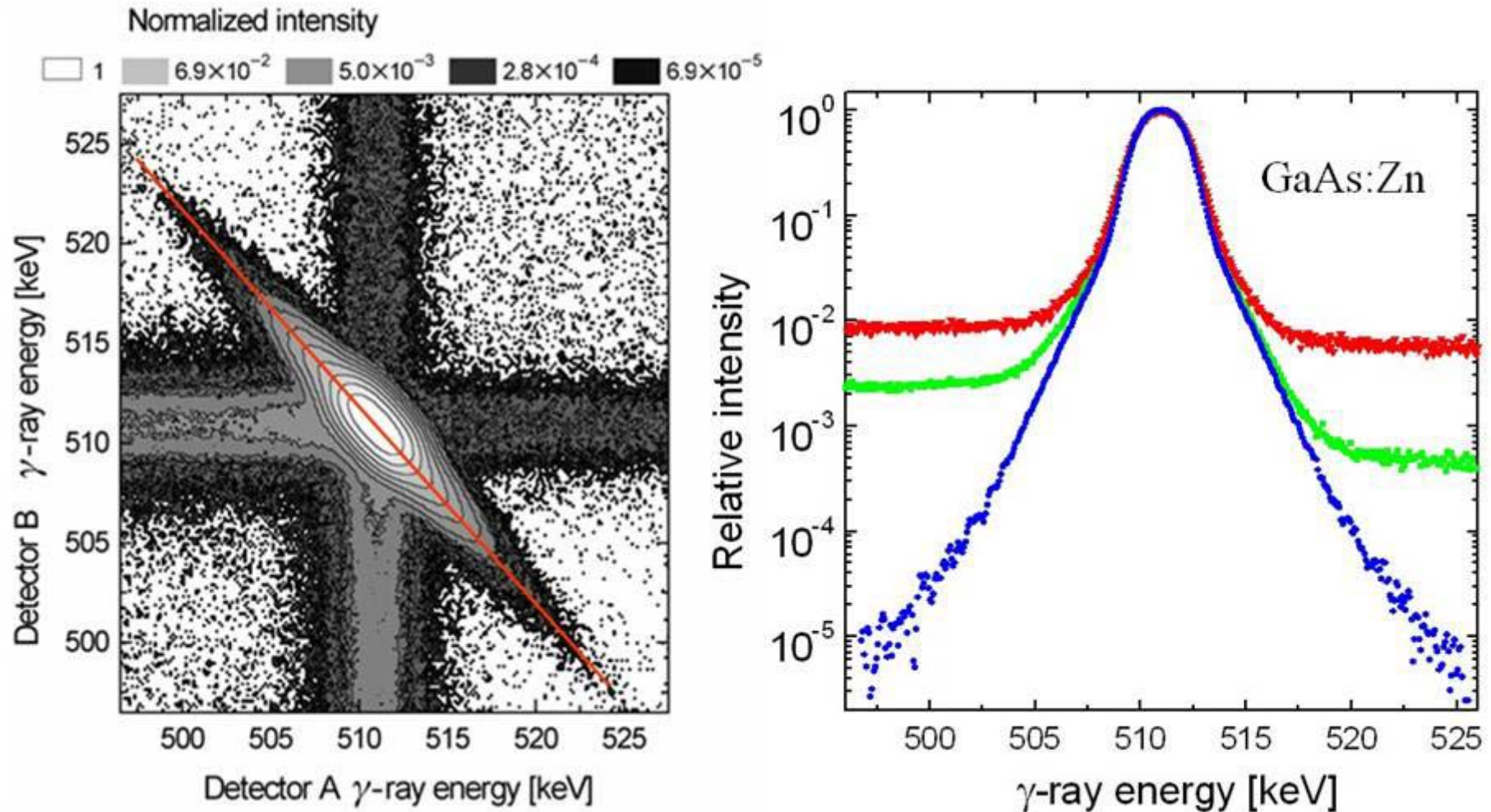


Doppler Coincidence Spectroscopy



- coincident detection of second annihilation γ reduces background
- use of a second Ge detector improves energy resolution of system

Doppler Coincidence Spectra

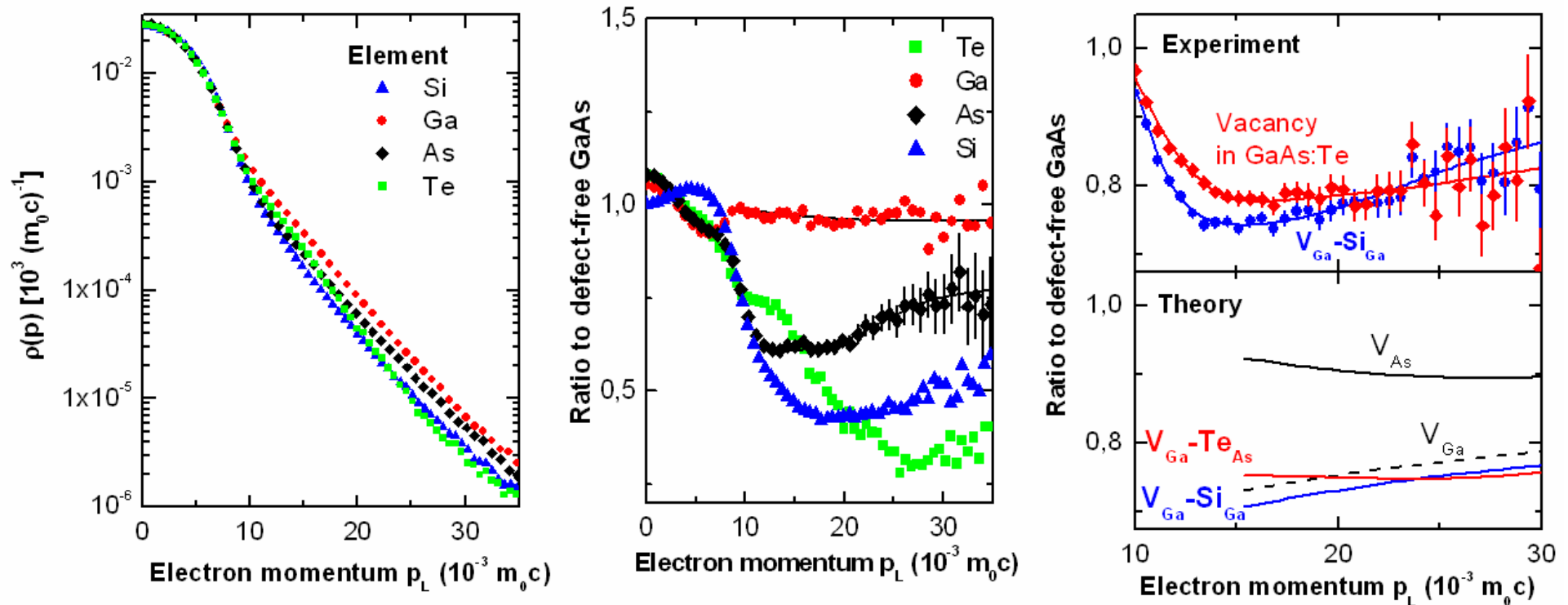


$$E_1 + E_2 = 2 m_0 c^2 = 1022 \text{ keV}$$



Doppler-Coincidence-Spectroscopy in GaAs

- Chemical sensitivity due to electrons at high momentum (core electrons)
- a single impurity atom aside a vacancy is detectable
- examples: $V_{Ga}-Te_{As}$ in GaAs:Te

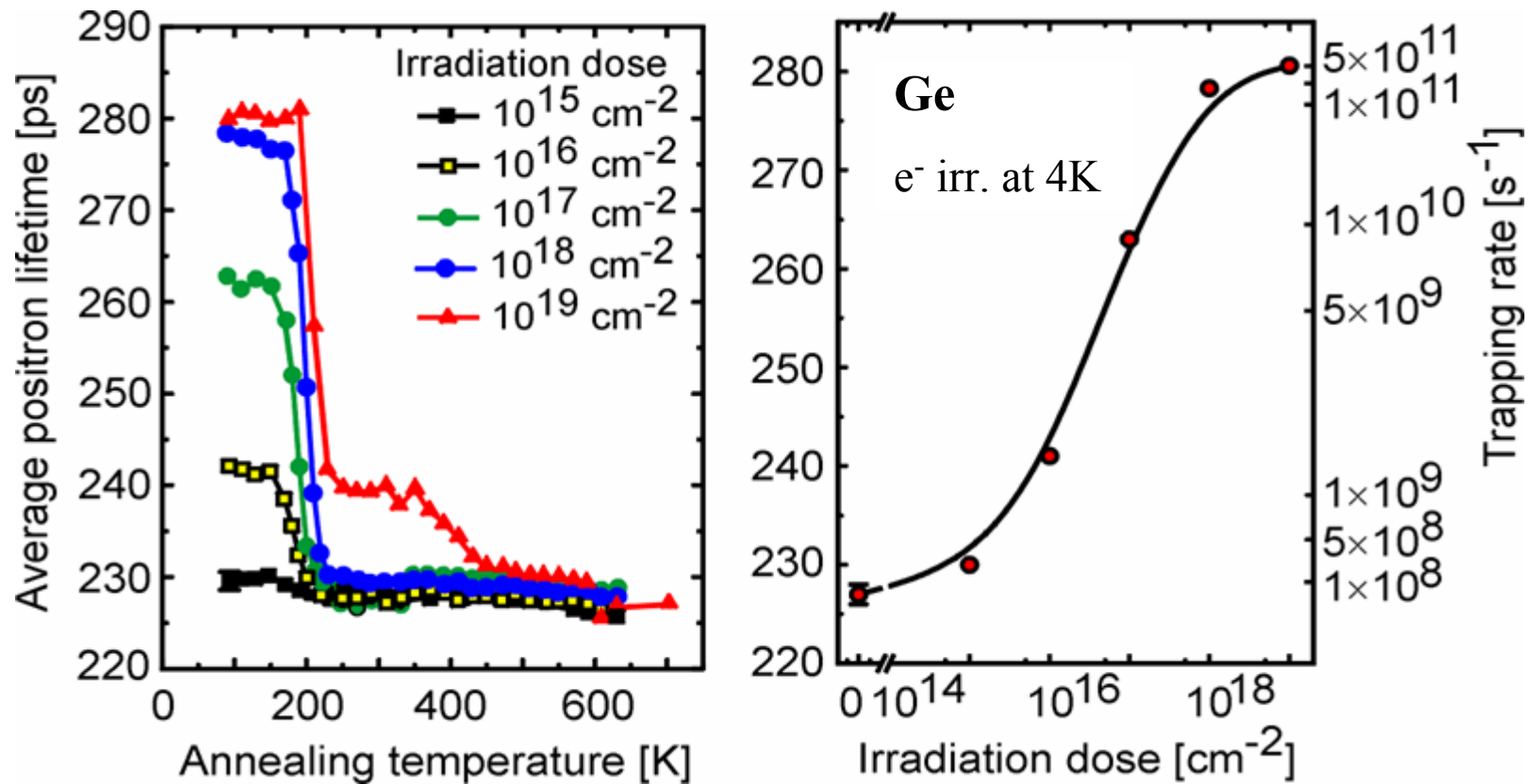


J. Gebauer et al., Phys. Rev. B **60** (1999) 1464



Defects in electron-irradiated Ge

- Electron irradiation (2 MeV @ 4K) induces Frenkel pairs (vacancy - interstitial pairs)
- steep annealing stage at 200 K
- at high irradiation dose: divacancies are formed (thermally more stable)

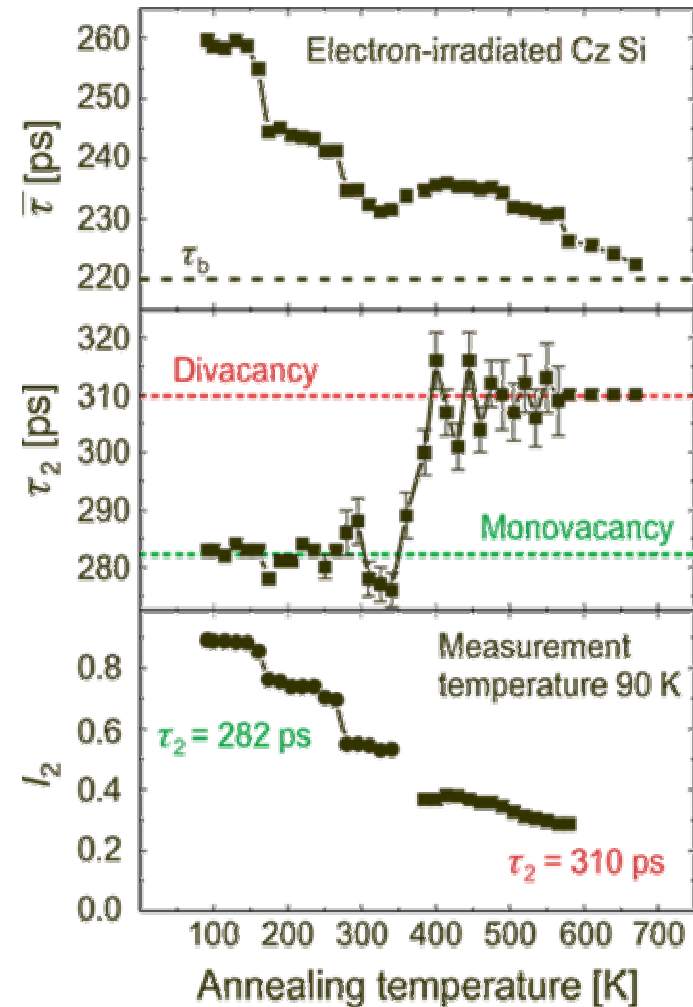
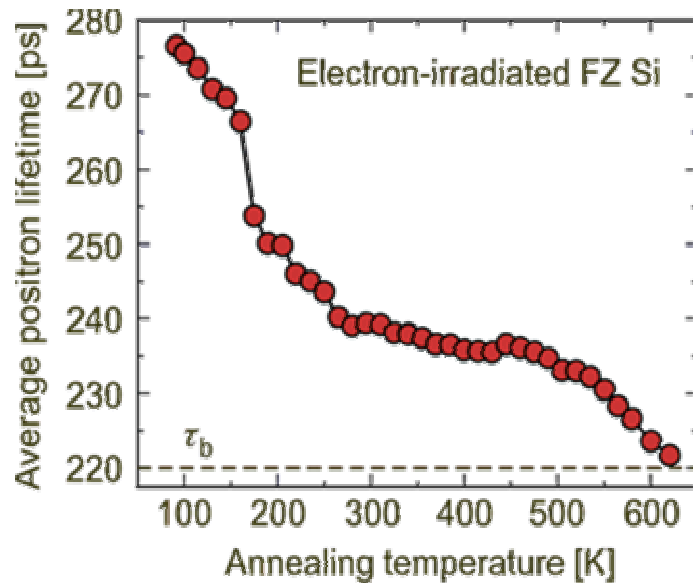


(Polity et al., 1997)



Low-temperature electron irradiation

- low-temperature electron irradiation was performed at 4K ($E_{e^-}=2$ MeV)
- annealing stage of monovacancies at about 170 K
- moving V_{Si} partly form divacancies
- divacancies anneal at about 550...650 K

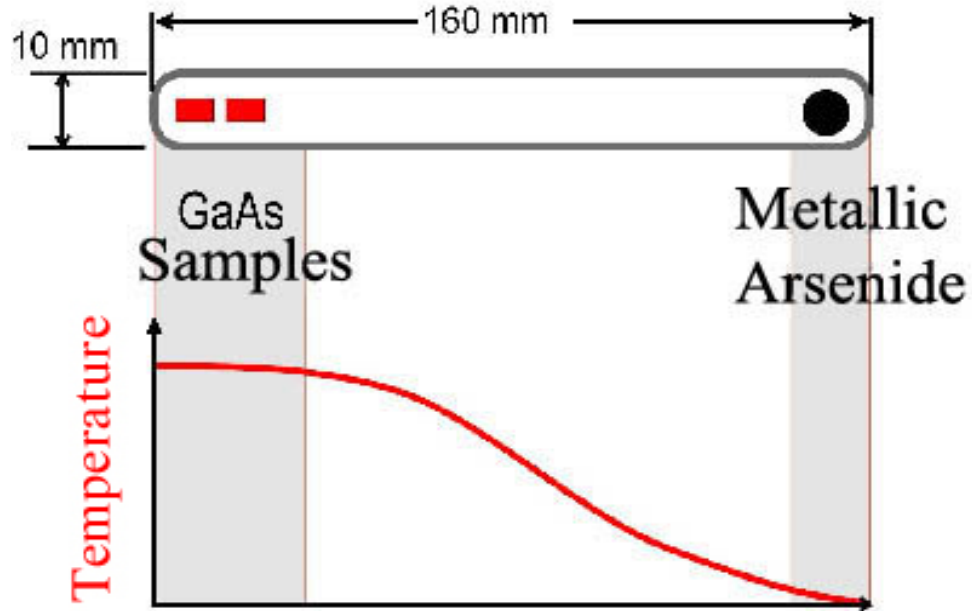


Polity et al., Phys. Rev. B **58** (1998) 10363



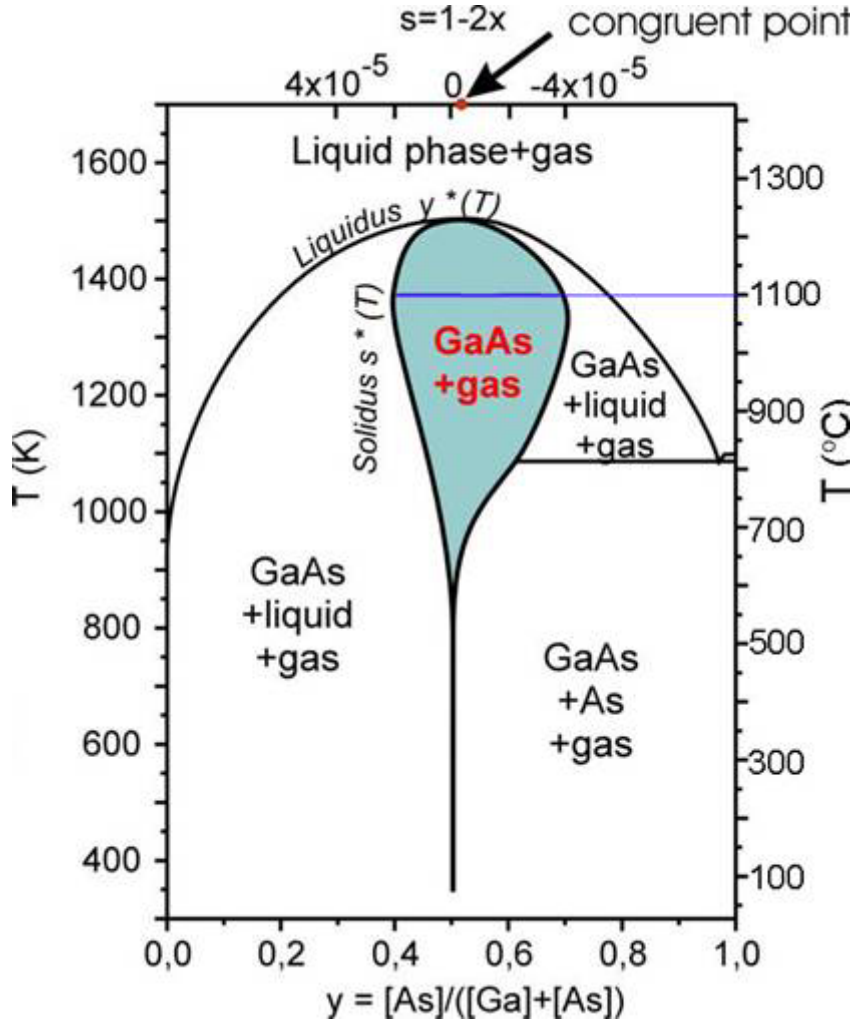
GaAs: annealing under defined As-partial pressure

- two-zone-furnace: Control of sample temperature **and** As partial pressure allows to navigate freely in phase diagram (existence area of compound)



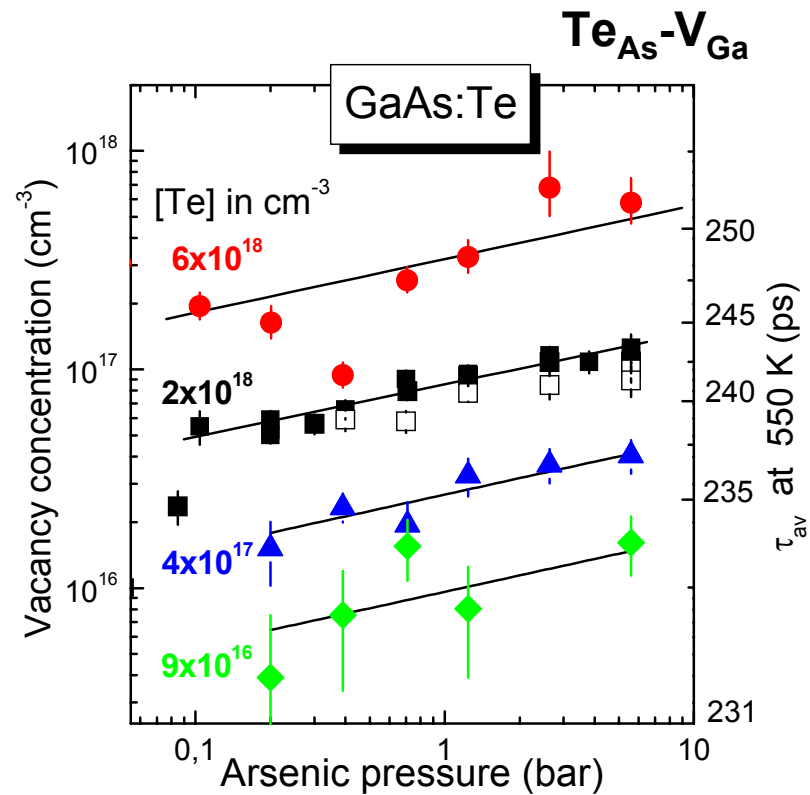
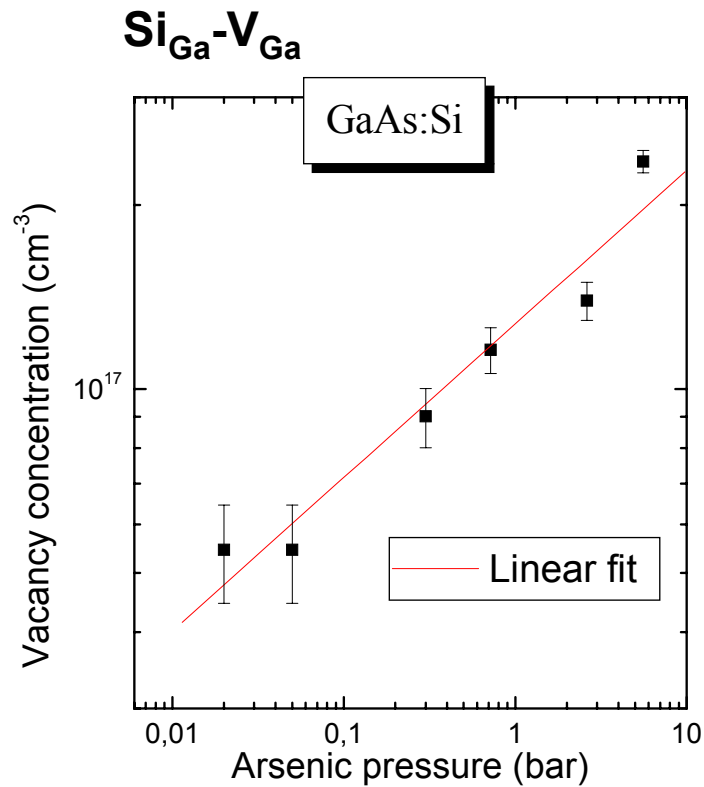
$T_{\text{sample}}: 1100^{\circ} \text{C}$

T_{As} : determines As-partial pressure



H. Wenzl et al., J. Cryst. Growth **109**, 191 (1991).

GaAs: Annealing under defined As pressure



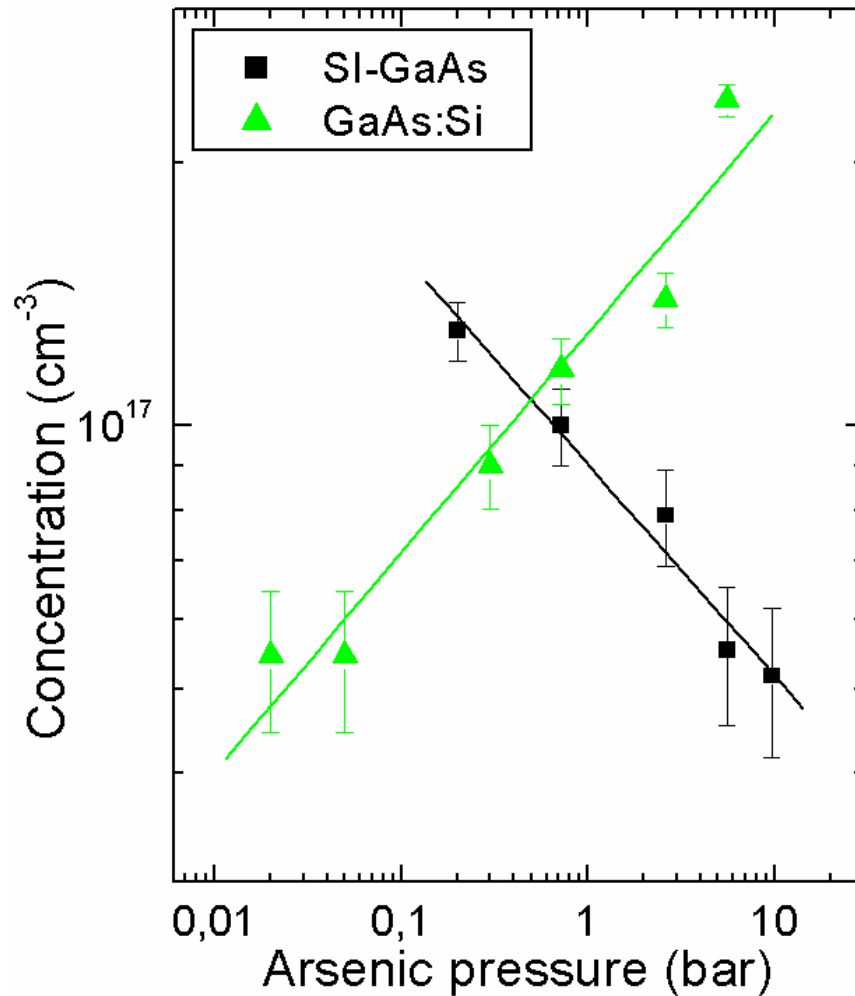
Thermodynamic reaction:
 $1/4 \text{As}_4^{\text{gas}} \leftrightarrow \text{As}_{\text{As}} + \text{V}_{\text{Ga}}$

Mass action law:
 $[\text{V}_{\text{Ga}}] = K_{\text{VG}} \times p_{\text{As}}^{1/4}$

Fit: $[\text{V}_{\text{Ga}}\text{-Dopant}] \sim p_{\text{As}}^n$
 $\rightarrow n = 1/4$

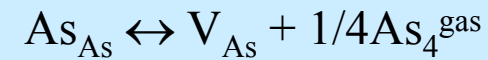
*J. Gebauer et al.,
 Physica B 273-274, 705 (1999)*

Comparison of doped and undoped GaAs



Bondarenko et al., 2003

Thermodynamic reaction:



Mass action law:

$$[\text{V}_{\text{As}}] = K_{\text{VAs}} \times p_{\text{As}}^{-1/4}$$

Fit: $[\text{V-complex}] \sim p_{\text{As}}^n$

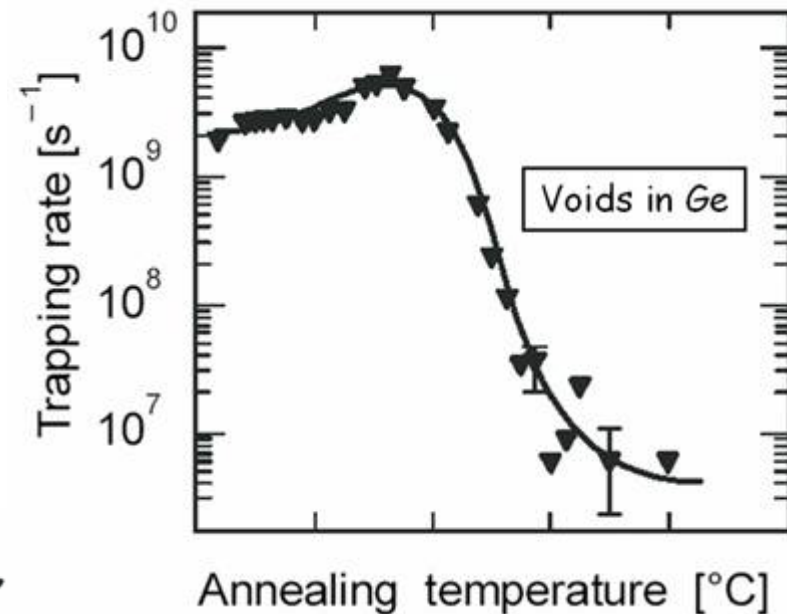
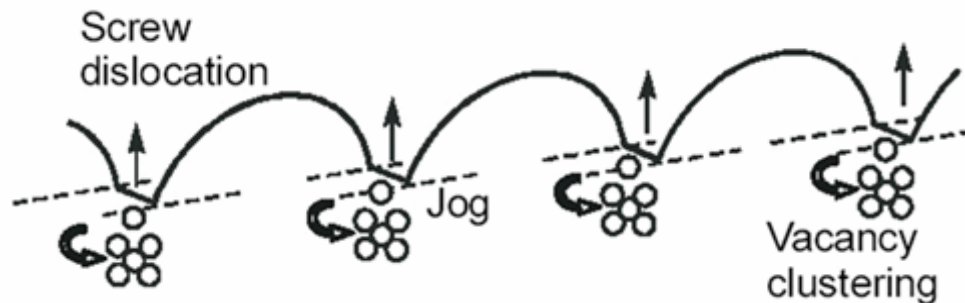
$$\rightarrow n = -1/4$$

As vacancy



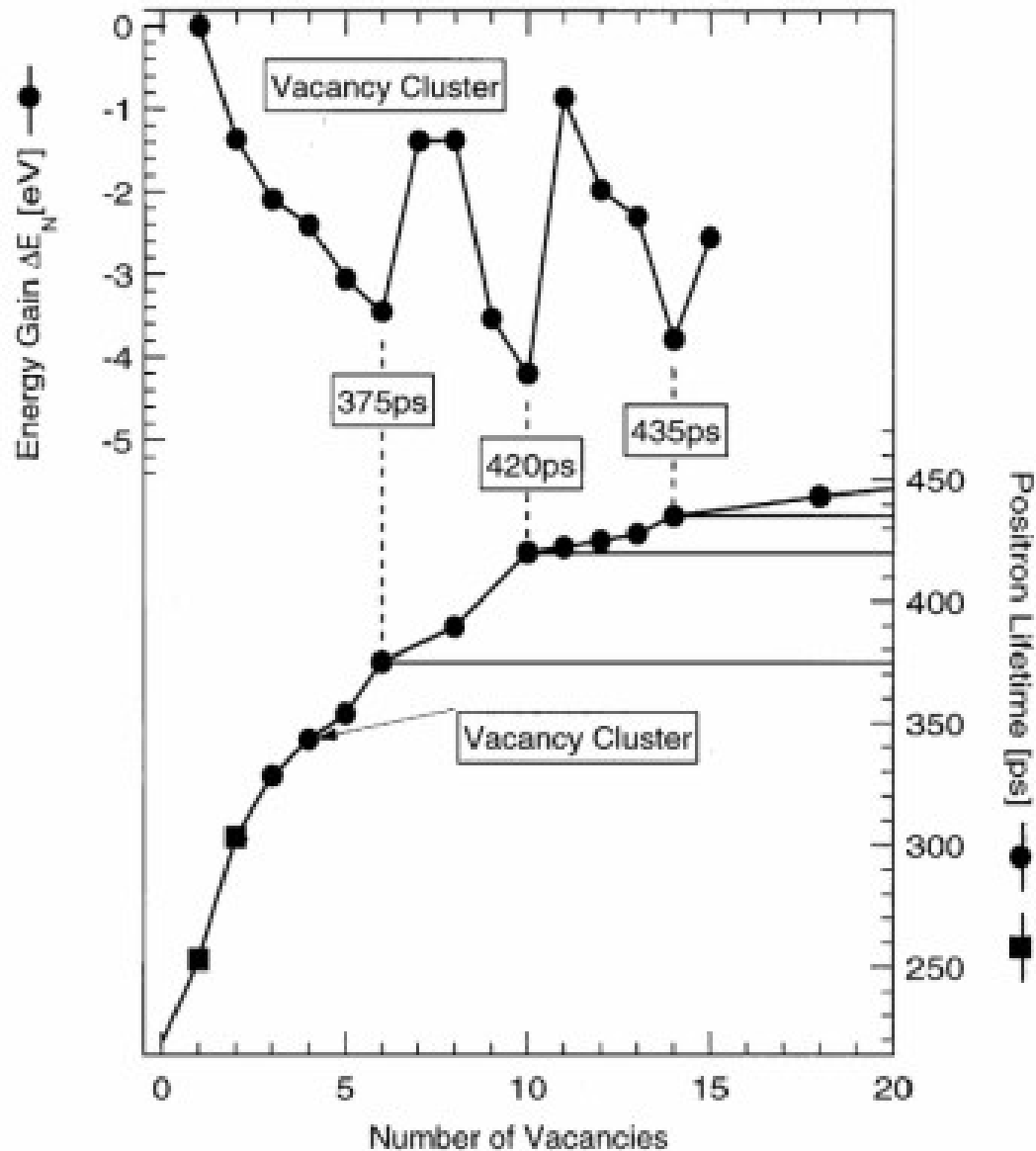
Vacancy clusters in semiconductors

- vacancy clusters were observed after neutron irradiation, ion implantation and plastic deformation
- due to large open volume (low electron density) \rightarrow positron lifetime increases distinctly
- example: plastically deformed Ge
- lifetime: $\tau = 525$ ps
- reason for void formation: jog dragging mechanism
- trapping rate of voids disappears during annealing experiment



Krause-Rehberg et al., 1993

Theoretical calculation of vacancy clusters in Si

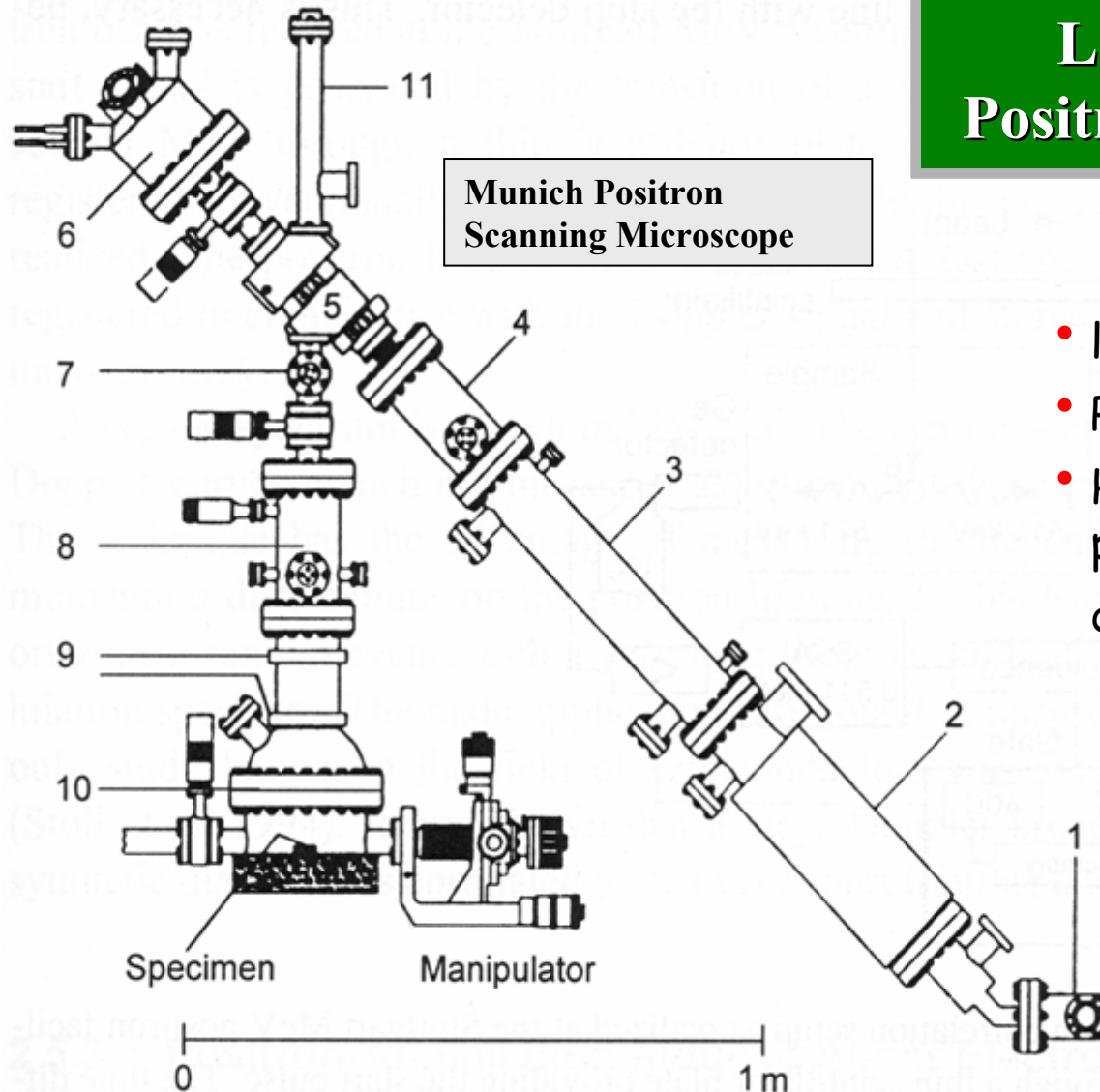


- there are cluster configurations with a large energy gain
- „Magic Numbers“ with 6, 10 und 14 vacancies
- positron lifetime increases distinctly with cluster size
- for $n > 10$ saturation effect, i.e. size cannot be determined

T.E.M. Staab et al.,
Physica B 273-274 (1999) 501-504



Lateral Resolution with Positron-Scanning-Microscope

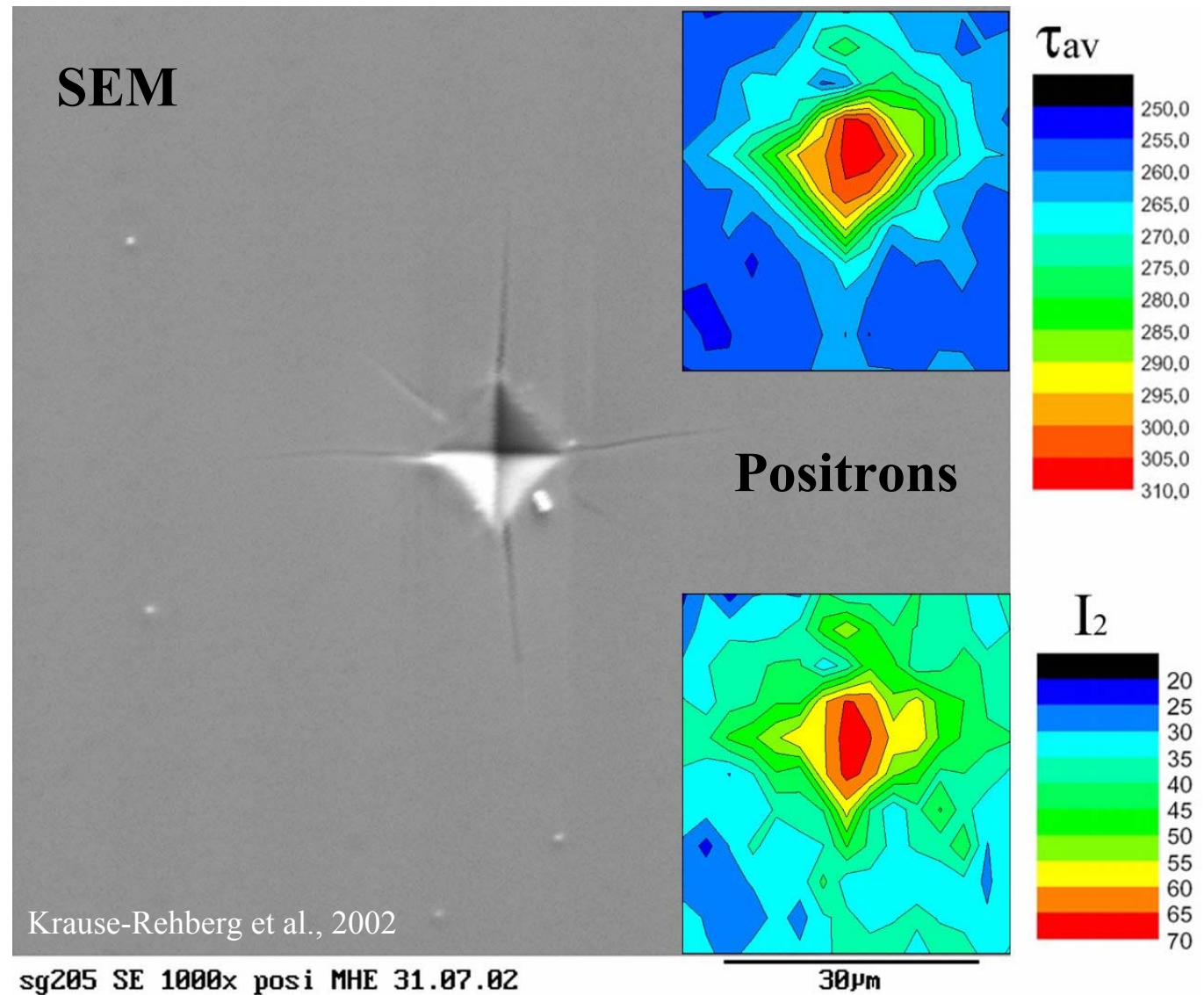


- lateral resolution $2 \mu\text{m}$
- Positron lifetime spectroscopy
- However lateral resolution principally limited by positron diffusion ($L_+ \approx 100\text{nm}$)

W. Triftshäuser et al., NIM B **130** (1997) 265

Microhardness indentation in GaAs

- Comparison of SEM and Munich Positron Scanning Microscope
- problem here at the moment: intensity
- hope: strong positron source at FRM-II Garching or EPOS project in Rossendorf

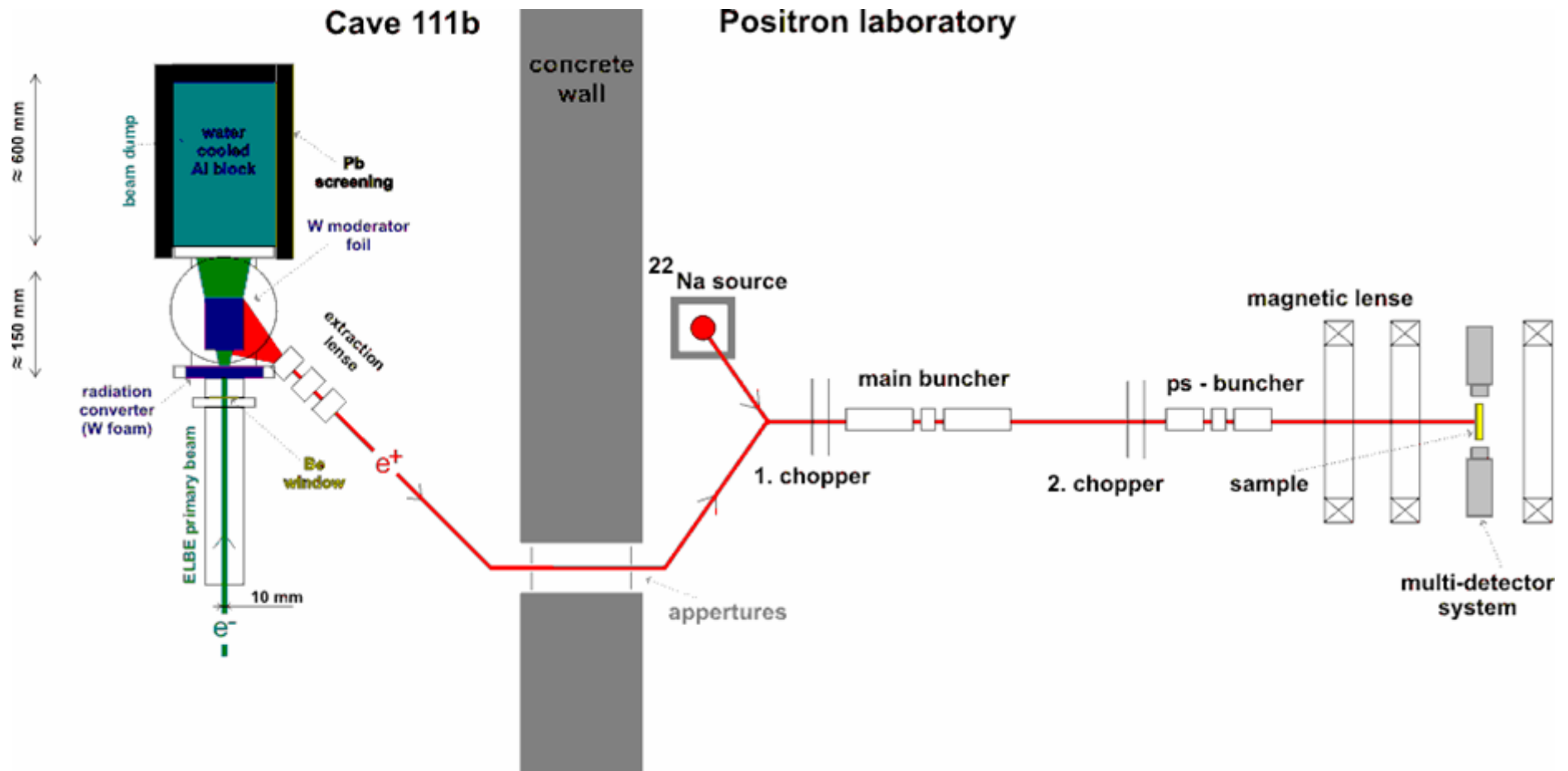


EPOS = ELBE Positron Source

- ELBE -> electron LINAC (40 MeV and up to 40 kW) in Research center Rossendorf (near Dresden)
- EPOS will be the combination of a positron lifetime spectrometer, Doppler coincidence, and AMOC
- User-dedicated facility
- main features:
 - ultra high-intensity bunched positron beam ($E_{\pm}=1...30$ keV)
 - very good time resolution by using the unique primary time structure of ELBE
 - high quality spectra by lifetime and Doppler spectroscopy in coincidence mode
 - fast lifetime mode (single detector mode) for kinetic investigations
 - very high count rate ($> 10^6$ s⁻¹) by multi-detector array
 - conventional source included for Doppler measurements (when primary beam is not available)
 - fully remote control via Internet by user



Schematic view of EPOS (ELBE Positron Source)



EPOS - Applications

Variety of applications in all field of materials science:

- defect-depth profiles due to surface modifications and ion implantation
- tribology (mechanical damage of surfaces)
- polymer physics (pores; interdiffusion; ...)
- low-k materials (thin high porous layers)
- defects in semiconductors, ceramics and metals
- epitaxial layers (growth defects, misfit defects at interface, ...)
- fast kinetics (e.g. precipitation processes in Al alloys; defect annealing; diffusion; ...)
- radiation resistance (e.g. space materials)
- many more ...



Conclusions

- Positrons are a unique tool for characterization of vacancy-type defects in solids
- Positrons are sensitive for charge state of vacancies in semiconductors
- vacancy clusters can easily be observed by positron lifetime spectroscopy (appear after irradiation and plastic deformation)

This presentation can be found as pdf-file on our Website:
<http://positron.physik.uni-halle.de>

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